

Stochastic subgrid turbulence parameterisation of eddy-eddy, eddy-topographic, eddy-meanfield and meanfield-meanfield interactions

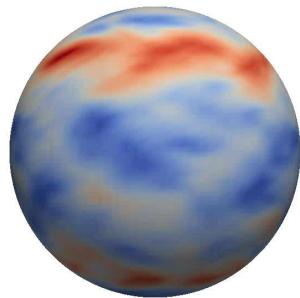
Vassili Kitsios & Jorgen S. Frederiksen

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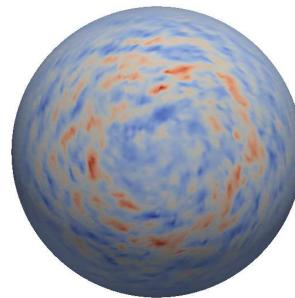


Motivation - Reduce GCM resolution dependence

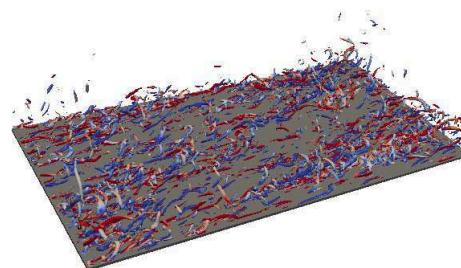
- Present work focussed on improving accuracy of GCMs, by reducing their resolution dependence.
- It is not possible to simulate all of the scales of motion, hence:
 - the large eddies are resolved by a computational grid
 - unresolved sub-grid scale (SGS) interactions are parameterised
- Typical approach: **Physical Hypothesis** → **Subgrid Model**
- Present approach: **Subgrid Model from DNS** → **Physical interpretation**
- Present stochastic subgrid modelling approach successfully applied to:



2-level QG Atmos.



2-level QG Ocean



300-level Boundary Layer

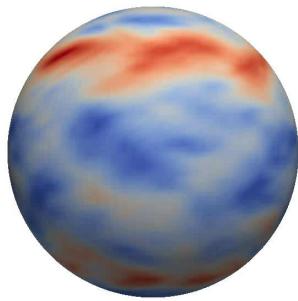
Kitsios et. al. (2012, JAS)

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Kitsios et. al. (2017, C&F)

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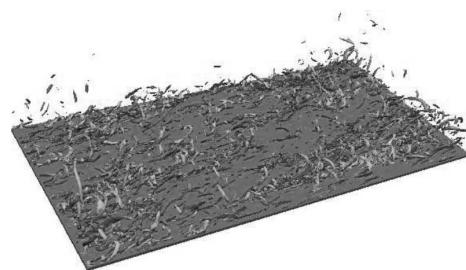
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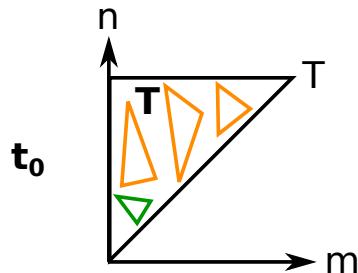
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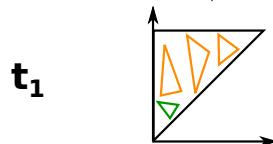
Decomposition of Scales: Triangular Truncation

DNS

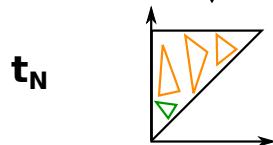
$$\frac{\partial \mathbf{q}}{\partial t} \equiv \dot{\mathbf{q}}_t = \mathcal{N}(\mathbf{q}; \mathbf{T})$$



timestep
 ↓



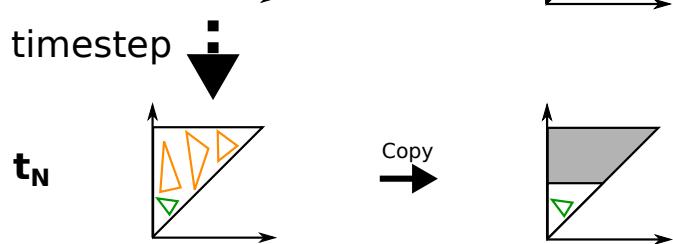
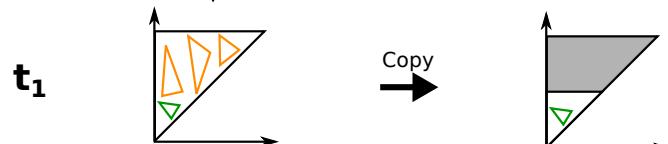
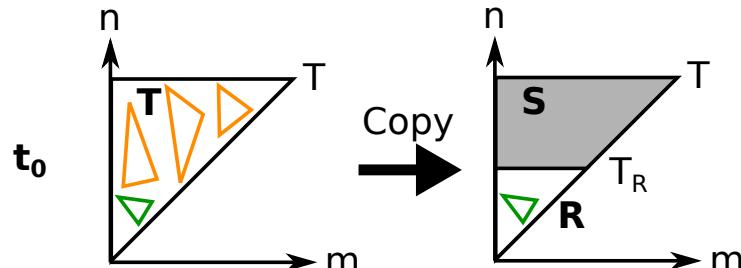
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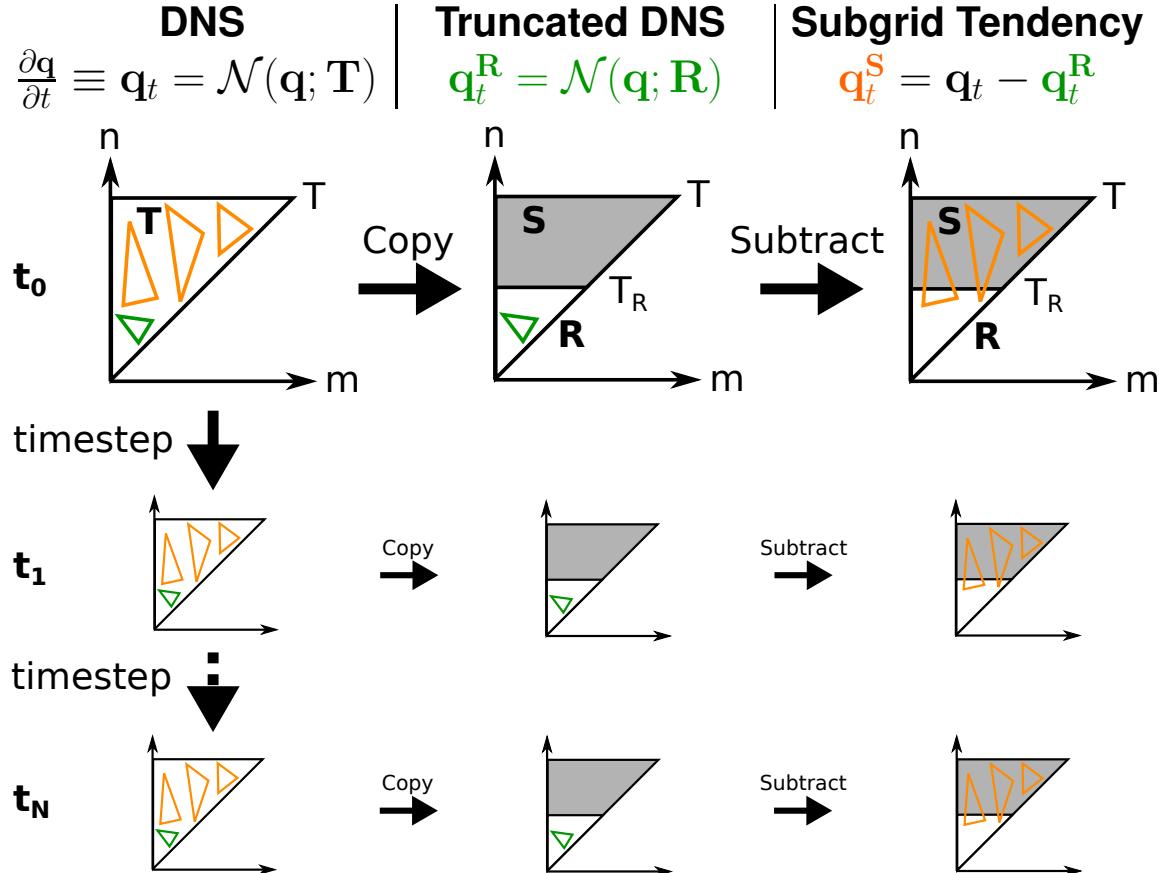
Decomposition of Scales: Triangular Truncation

$$\text{DNS} \quad \frac{\partial \mathbf{q}}{\partial t} \equiv \mathbf{q}_t = \mathcal{N}(\mathbf{q}; \mathbf{T})$$

$$\text{Truncated DNS} \quad \mathbf{q}_t^R = \mathcal{N}(\mathbf{q}; \mathbf{R})$$



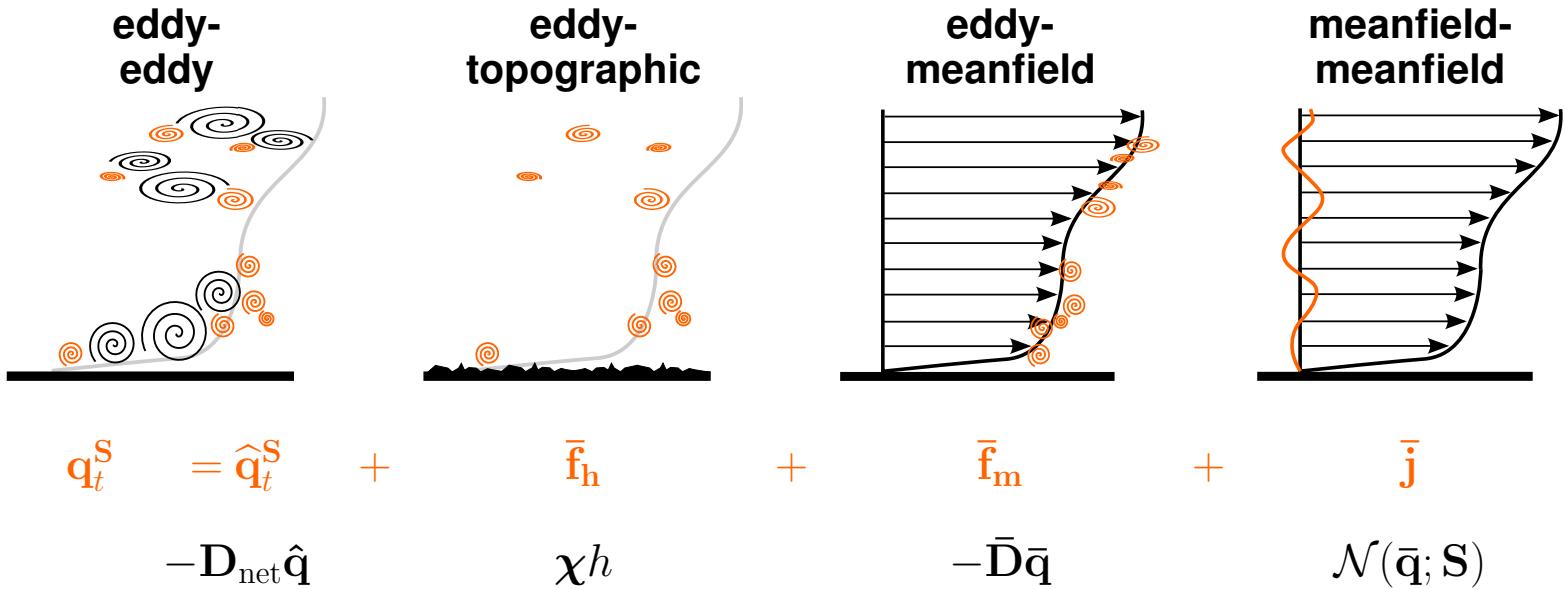
Decomposition of Scales: Triangular Truncation



- In most primitive form, subgrid modelling relates \mathbf{q}_t^S to \mathbf{q} .

Subgrid interactions

- **Eddy-eddy:** *subgrid eddies* and *resolved eddies*.
- **Eddy-topographic:** *subgrid eddies* and *resolved topography*.
- **Eddy-meanfield:** *subgrid eddies* and *resolved meanfield*.
- **Meanfield-meanfield:** *subgrid meanfield* and *resolved meanfield*.



- Functional forms derived from closure theory (Frederiksen, 1999, JAS)

Eddy-Eddy Dissipation

- Subgrid interactions are *local* in wavenumbers space, not grid space.
- Unique Matrix captures level interactions per scale (Frederiksen & Kepert, 2006)

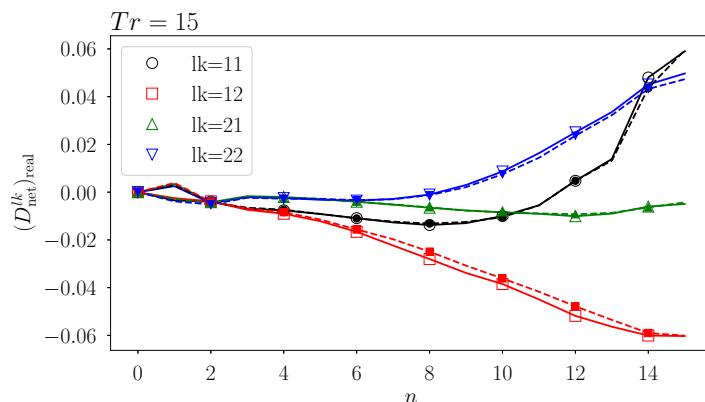
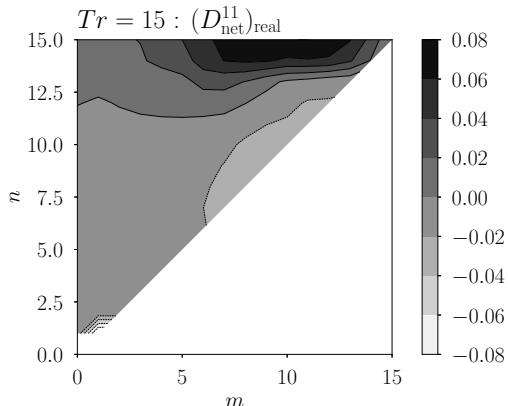
Deterministic

$$\begin{aligned}\hat{\mathbf{q}}_t^S(t) &= -D_{\text{net}} \hat{\mathbf{q}}(t) \\ \langle \hat{\mathbf{q}}_t^S(t) \hat{\mathbf{q}}^\dagger(t) \rangle &= -D_{\text{net}} \langle \hat{\mathbf{q}}(t) \hat{\mathbf{q}}^\dagger(t) \rangle \\ D_{\text{net}} &= -\langle \hat{\mathbf{q}}_t^S(t) \hat{\mathbf{q}}^\dagger(t) \rangle \langle \hat{\mathbf{q}}(t) \hat{\mathbf{q}}^\dagger(t) \rangle^{-1}\end{aligned}$$

Stochastic

$$\begin{aligned}\hat{\mathbf{q}}_t^S(t) &= -D_d \hat{\mathbf{q}}(t) + \hat{\mathbf{f}}(t) \\ D_d \text{ incorporates history} \\ D_b &\propto \text{variance of } \hat{\mathbf{f}}(t)\end{aligned}$$

- At $T_R = 15$ baroclinic instability is not resolved
- Dissipation anisotropic, negative for many scales, off diagonals significant

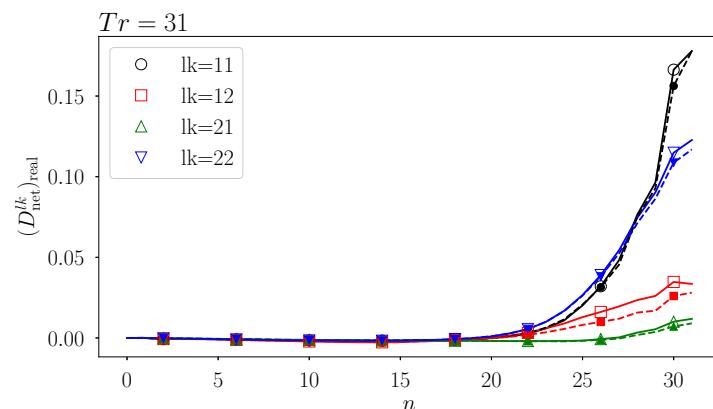
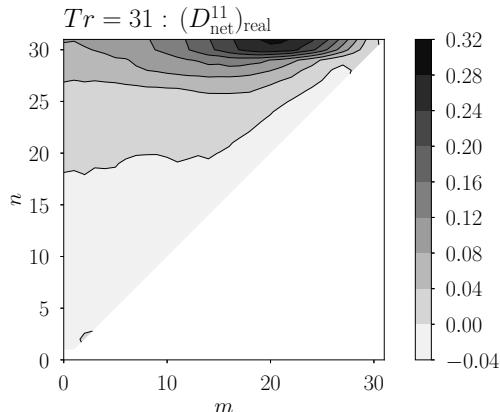


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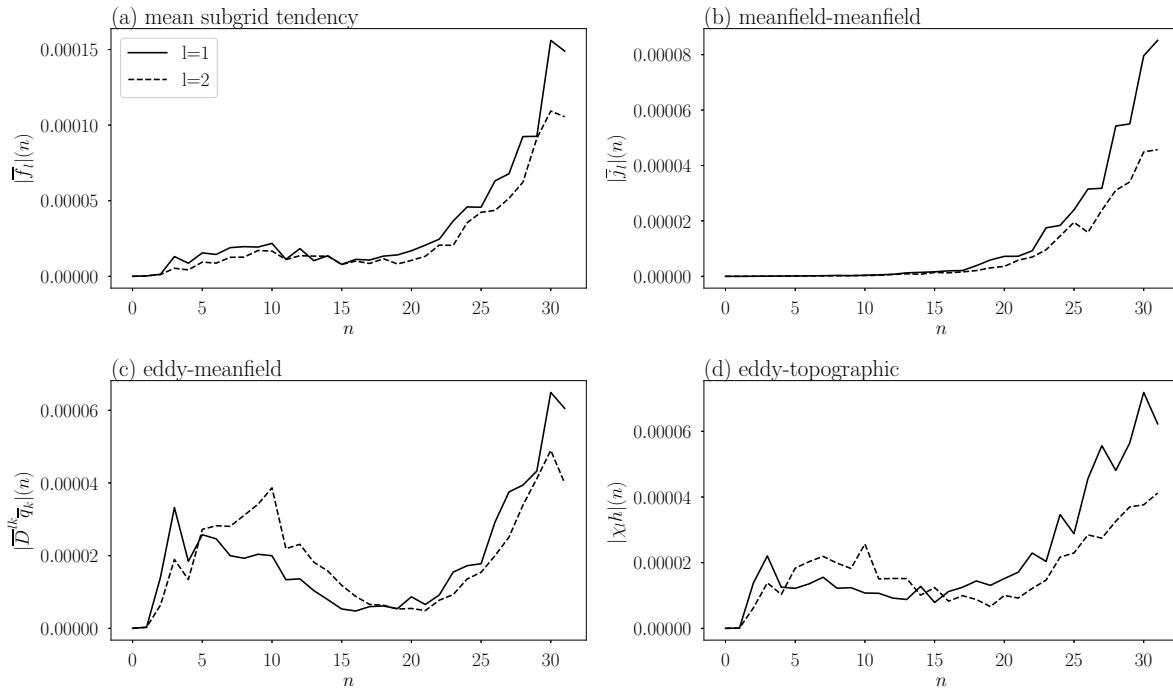
Deterministic	Stochastic
$\hat{\mathbf{q}}_t^S(t) = -D_{\text{net}} \hat{\mathbf{q}}(t)$ $\langle \hat{\mathbf{q}}_t^S(t) \hat{\mathbf{q}}^\dagger(t) \rangle = -D_{\text{net}} \langle \hat{\mathbf{q}}(t) \hat{\mathbf{q}}^\dagger(t) \rangle$ $D_{\text{net}} = -\langle \hat{\mathbf{q}}_t^S(t) \hat{\mathbf{q}}^\dagger(t) \rangle \langle \hat{\mathbf{q}}(t) \hat{\mathbf{q}}^\dagger(t) \rangle^{-1}$	$\hat{\mathbf{q}}_t^S(t) = -D_d \hat{\mathbf{q}}(t) + \hat{\mathbf{f}}(t)$ D_d incorporates history $D_b \propto$ variance of $\hat{\mathbf{f}}(t)$

- At $T_R = 31$ baroclinic instability is better resolved
- Dissipation isotropic, negative for fewer scales, off diagonals smaller



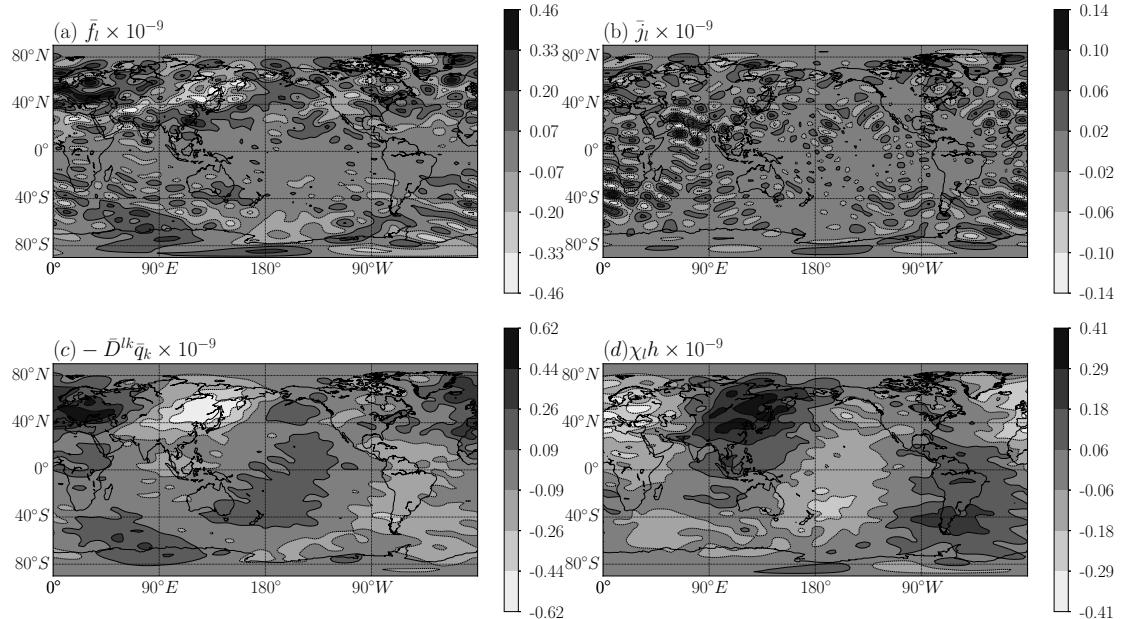
Decompose Mean Subgrid Tendency $\equiv \overline{\mathbf{q}_t^S} \equiv \bar{\mathbf{f}}_h + \bar{\mathbf{f}}_m + \bar{\mathbf{j}}$

- **meanfield-meanfield:** evaluate $\bar{\mathbf{j}} = \mathcal{N}(\bar{\mathbf{q}}; \mathbf{S}) \equiv \mathcal{N}(\bar{\mathbf{q}}; \mathbf{T}) - \mathcal{N}(\bar{\mathbf{q}}; \mathbf{R})$
- **eddy-meanfield:** $\bar{\mathbf{f}}_m = -\bar{\mathbf{D}}\bar{\mathbf{q}}$
determine $\bar{\mathbf{D}}$ via regression of $[\overline{\mathbf{q}_t^S} - \bar{\mathbf{j}}]$ with $\bar{\mathbf{q}}$ over 500 climate states
- **eddy-topographic:** determined by the remainder $\bar{\mathbf{f}}_h = \overline{\mathbf{q}_t^S} - \bar{\mathbf{j}} + \bar{\mathbf{D}}\bar{\mathbf{q}}$



Decompose Mean Subgrid Tendency $\equiv \overline{\mathbf{q}_t^S} \equiv \bar{\mathbf{f}}_h + \bar{\mathbf{f}}_m + \bar{\mathbf{j}}$

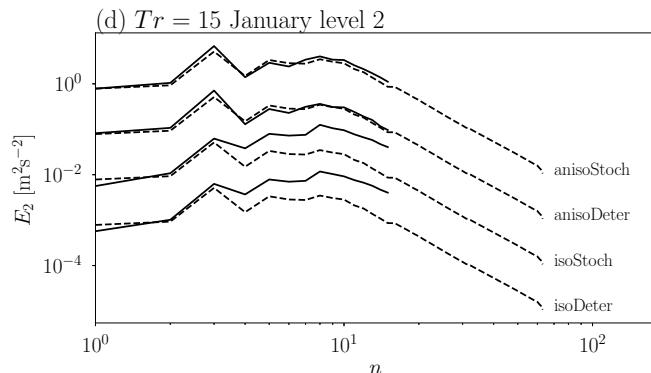
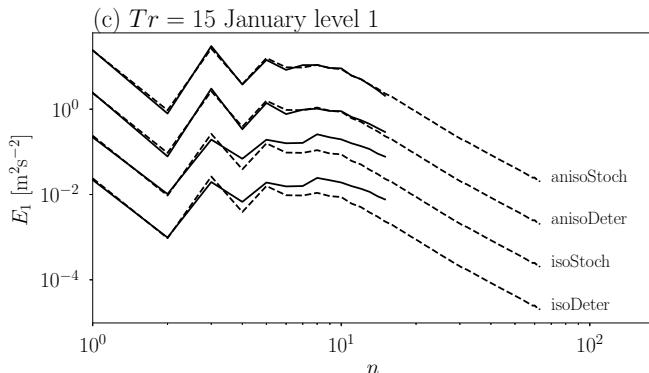
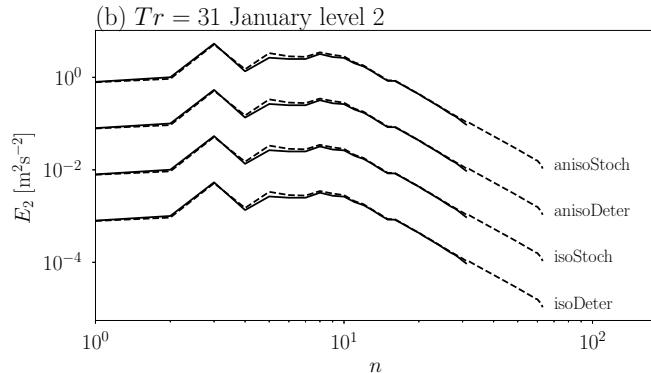
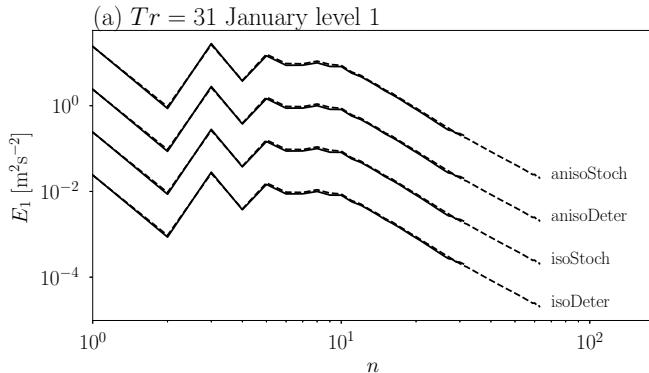
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- Similar to previous evaluation of closure terms (O'Kane & Frederiksen, 2008)

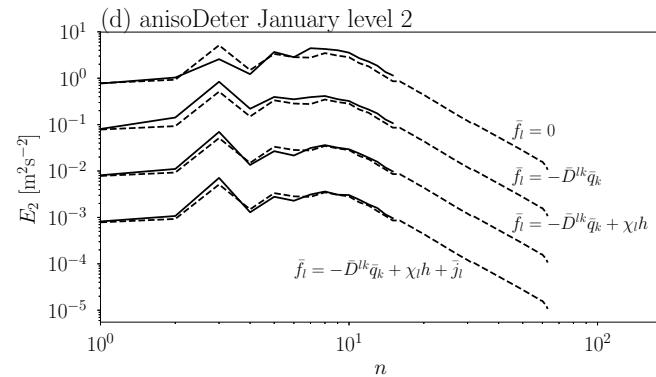
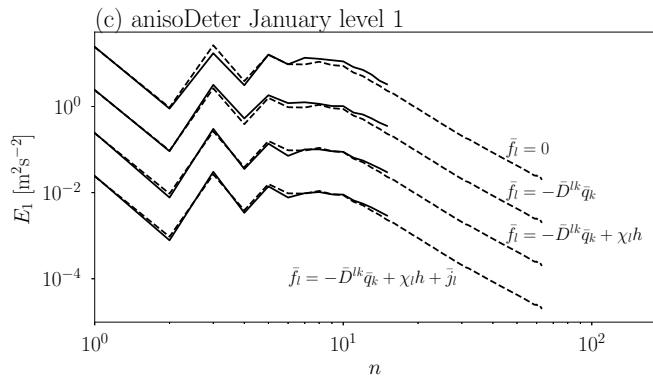
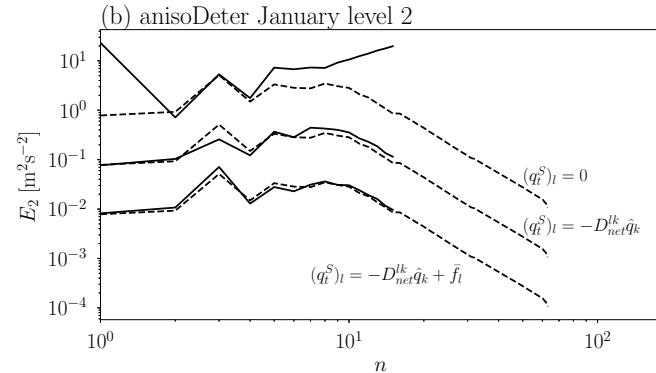
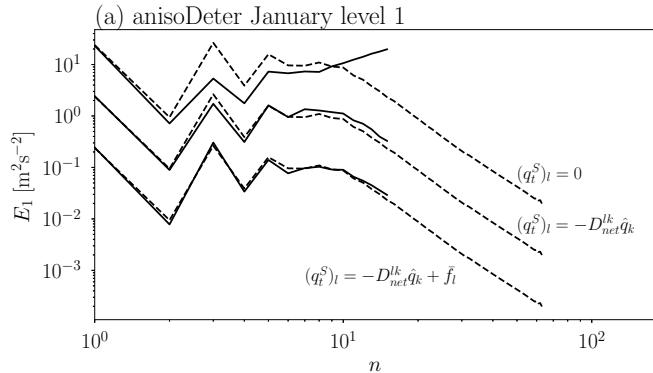
LES : Eddy-Eddy Variants ; Full Mean Subgrid Tendency

- For $T_R = 31$ **stochastic** and **deterministic** variants with **anisotropic** and **isotropic** eddy-eddy coefficients reproduce the DNS for all scales
- For $T_R = 15$ **anisotropic** coefficients are required



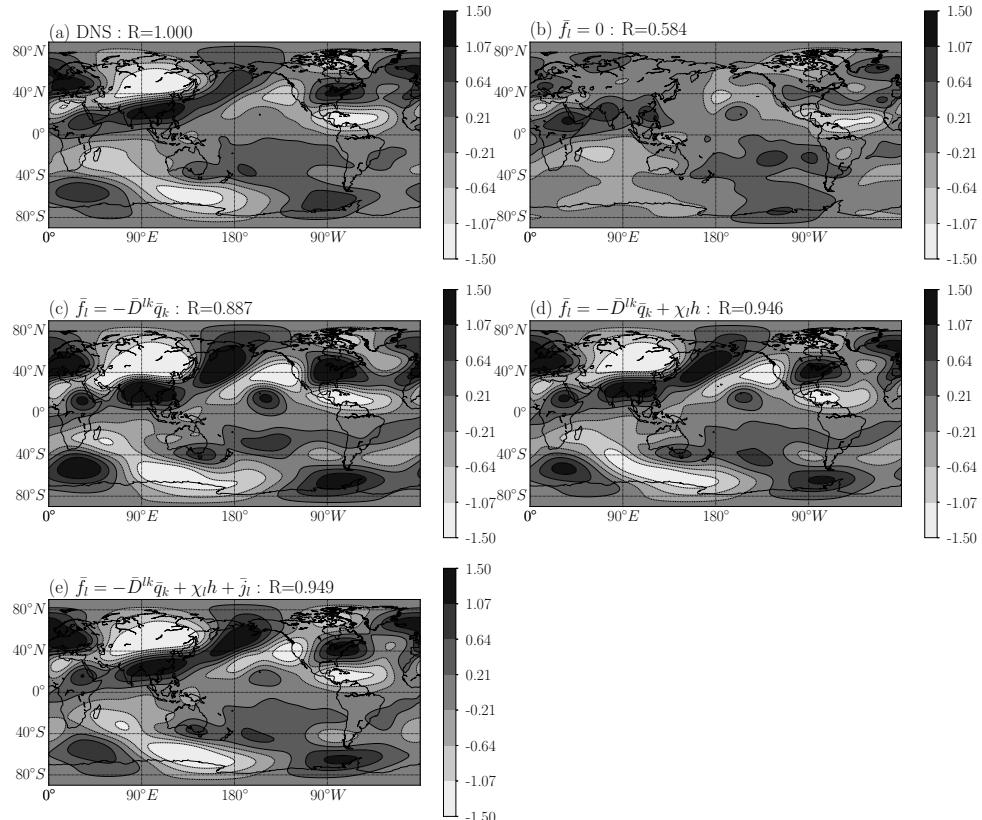
Anisotropic Deterministic LES at $T_R = 15$

- Eddy-eddy parameterisation has dominant impact
- Best agreement requires parameterisation of all subgrid interactions



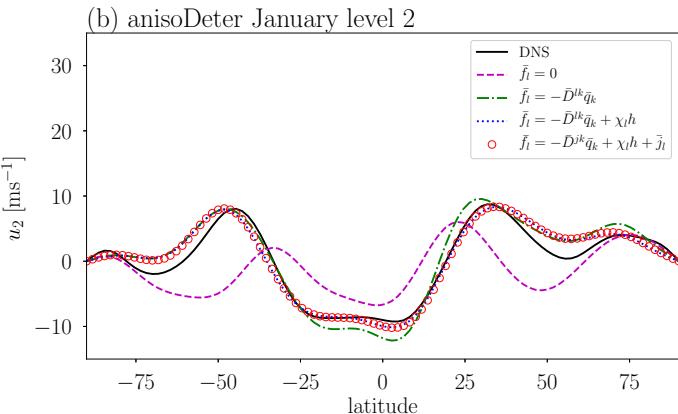
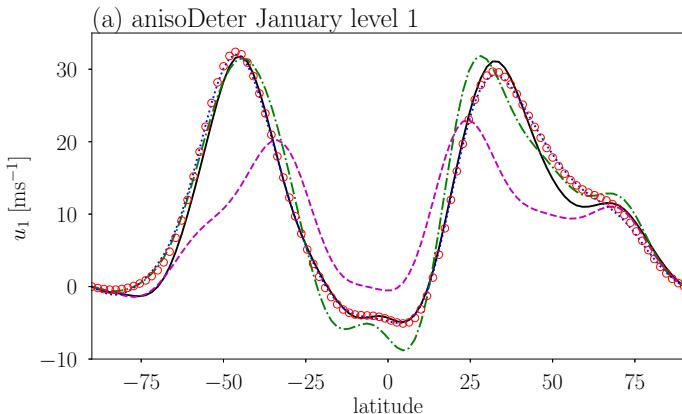
Anisotropic Deterministic LES at $T_R = 15$

- Likewise best pattern correlation of mean nonzonal streamfunction requires parameterisation of all subgrid interactions



Anisotropic Deterministic LES at $T_R = 15$

- Likewise best pattern correlation of mean nonzonal streamfunction requires parameterisation of all subgrid interactions
- Also true to for mean zonal velocity

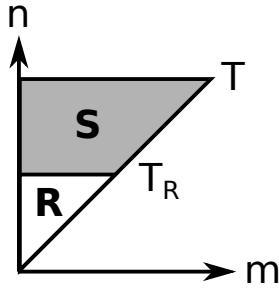


Concluding Remarks

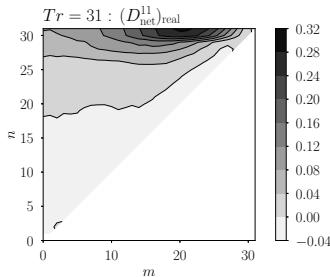
- Stochastic subgrid modelling approach used to determine parameterisation coefficients representing the interactions between the:
 - subgrid eddies and resolved eddies (**eddy-eddy**)
 - subgrid eddies and resolved topography (**eddy-topographic**)
 - subgrid eddies and resolved meanfield (**eddy-meanfield**)
 - subgrid meanfield and resolved meanfield (**meanfield-meanfield**)
- Note there are no tuning parameters in this approach.
- The eddy-eddy parameterisation has the dominant contribution to reproducing the DNS spectra across all scales.
- However, parameterisation of **all subgrid interactions** required for best agreement of the DNS and LES in terms of:
 - kinetic energy spectra
 - mean nonzonal streamfunction
 - mean zonal velocity

Questions

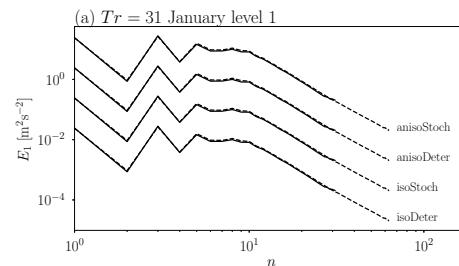
1) Decompose DNS



2) Calculate Dissipation



3) Verify with LES



- **Kitsios, V., Sillero, J.A., Frederiksen, J.S. & Soria, J., 2017**, Scale and Reynolds number dependence of stochastic subgrid energy transfer in turbulent channel flow, Computers and Fluids, Vol. 151, pp 132-143.
- **Kitsios, V., Frederiksen, J.S. & Zidikheri, M.J., 2013**, Scaling laws for parameterisations of subgrid eddy-eddy interactions in simulations of oceanic circulations, Ocean Modelling, 68, pp 88-105.
- **Kitsios, V., Frederiksen, J.S. & Zidikheri, M.J., 2012**, Subgrid model with scaling laws for atmospheric simulations, Journal of the Atmospheric Sciences, 69, pp 1427-1445.
- **O’Kane, T.J. & Frederiksen, J.S., 2008**, Statistical dynamical subgrid-scale parameterizations for geophysical flows, Phys. Scr., T132, 014033, (11pp).
- **Frederiksen & Kepert, J.S., 2006**, Dynamical subgrid-scale parameterizations from direct numerical simulations, Journal of the Atmospheric Sciences, 63, pp 3006-3019.
- **Frederiksen, J.S., 1999**, Subgrid-scale parameterizations of eddy-topographic force, eddy viscosity and stochastic backscatter for flow over topography, Journal of the Atmospheric Sciences, 56, pp 1481-1494.

Thank You

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